

Enthalpy $H(T,p)$ $H \equiv U + pV$

Chemical reactions and biological processes usually take place under constant pressure and with reversible pV work. Enthalpy turns out to be an especially useful function of state under those conditions.

$$\begin{array}{ccc} \text{gas } (p, T_1, V_1) & \stackrel{\text{reversible}}{=} & \text{gas } (p, T_2, V_2) \\ & \text{const. } p & \\ U_1 & & U_2 \end{array}$$

$$\Delta U = q + w = q_p - p\Delta V$$

$$\Delta U + p\Delta V = q_p$$

$$\Delta U + \Delta(pV) = q_p \Rightarrow \Delta(U + pV) = q_p$$

define as H

$$H \equiv U + pV \Rightarrow \Delta H = q_p \quad \text{for a reversible constant } p \text{ process}$$

Choose $H(T,p) \Rightarrow dH = \left(\frac{\partial H}{\partial T}\right)_p dT + \left(\frac{\partial H}{\partial p}\right)_T dp$

What are $\left(\frac{\partial H}{\partial T}\right)_p$ and $\left(\frac{\partial H}{\partial p}\right)_T$?

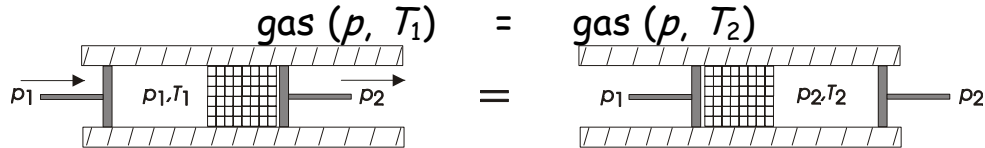
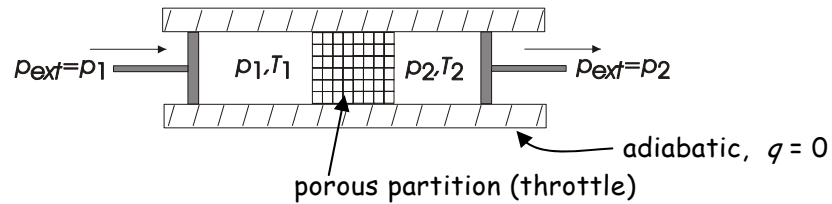
- $\left(\frac{\partial H}{\partial T}\right)_p \Rightarrow$ for a reversible process at constant p ($dp = 0$)

$$dH = \delta q_p \quad \text{and} \quad dH = \left(\frac{\partial H}{\partial T}\right)_p dT$$

$$\Rightarrow \delta q_p = \left(\frac{\partial H}{\partial T}\right)_p dT \quad \text{but} \quad \delta q_p = C_p dT \quad \text{also}$$

$$\therefore \boxed{\left(\frac{\partial H}{\partial T}\right)_p = C_p}$$

- $\left(\frac{\partial H}{\partial p}\right)_T \Rightarrow$ Joule-Thomson expansion



$$w = p_1 V_1 - p_2 V_2 \Rightarrow \Delta U = q + w = p_1 V_1 - p_2 V_2 = -\Delta(pV)$$

$$\therefore \Delta U + \Delta(pV) = 0 \Rightarrow \Delta(U + pV) = 0$$

$$\therefore \boxed{\Delta H = 0}$$

Joule-Thomson is a constant Enthalpy process.

$$dH = C_p dT + \left(\frac{\partial H}{\partial p}\right)_T dp \Rightarrow C_p dT = -\left(\frac{\partial H}{\partial p}\right)_T dp_H$$

$$\Rightarrow \left(\frac{\partial H}{\partial p}\right)_T = -C_p \left(\frac{\partial T}{\partial p}\right)_H \leftarrow \text{can measure this } \left(\frac{\Delta T}{\Delta p}\right)_H$$

Define $\lim_{\Delta p \rightarrow 0} \left(\frac{\Delta T}{\Delta p}\right)_H = \left(\frac{\partial T}{\partial p}\right)_H \equiv \mu_{JT} \leftarrow$ Joule-Thomson Coefficient

$$\therefore \boxed{\left(\frac{\partial H}{\partial p}\right)_T = -C_p \mu_{JT}} \quad \text{and} \quad \boxed{dH = C_p dT - C_p \mu_{JT} dp}$$

Proof that $\bar{C}_p = \bar{C}_v + R$ for an ideal gas

$$\begin{aligned} \bar{C}_p &= \left(\frac{\partial \bar{H}}{\partial T} \right)_p, & \bar{C}_v &= \left(\frac{\partial \bar{U}}{\partial T} \right)_v \\ \bar{H} &= \bar{U} + p\bar{V}, & p\bar{V} &= RT \\ \left(\frac{\partial \bar{H}}{\partial T} \right)_p &= \left(\frac{\partial \bar{U}}{\partial T} \right)_p + p \left(\frac{\partial \bar{V}}{\partial T} \right)_p \\ \bar{C}_p &= \bar{C}_v + \underbrace{\left(\frac{\partial \bar{U}}{\partial \bar{V}} \right)_T}_{=0 \text{ for ideal gas}} \left(\frac{\partial \bar{V}}{\partial T} \right)_p + \cancel{p} \left(\frac{R}{\cancel{p}} \right) \\ \therefore \bar{C}_p &= \bar{C}_v + R \end{aligned}$$